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# Similitude in Hydraulic Modelling

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#### **Background Hydraulic Modelling**











Clockwise from to left: Ski jump at Karakaya hydropower dam, sediment transport, tidal turbine (SeaGen), weir, flooding, breaking wave, floating wave energy converter (Plemais)





# Similarity: Gebidem Dam, Switzerland

#### Laboratory model at 1:30

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Jet trajectory  $\checkmark$  Air concentration **\stackrel{\scriptstyle{\leftarrow}}{\phantom{\phantom{\leftarrow}}}** 

Full scale (prototype)



Scale factor  $\lambda = L_P/L_M$  with  $L_P$  = characteristic length in the **P**rototype and  $L_M$  = corresponding length in the **M**odel 4

A physical scale model satisfying *mechanical similarity* is completely similar to its prototype and involves no scale effects (Hughes 1993).

Mechanical similarity requires three criteria to be satisfied:

#### (i) Geometric similarity:

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Similarity in shape, i.e. all length dimensions in the model are  $\lambda$  times shorter than of its prototype ( $\lambda = L_P/L_M$ )

#### (ii) Kinematic similarity:

Geometric similarity and similarity of motion between model and prototype particles

#### (iii) **Dynamic similarity**:

Requires geometric and kinematic similarity and in addition that all force ratios in the two systems are identical





#### The **most relevant forces** in fluid dynamics are:

• Inertial force

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- Gravitational force
- Viscous force

- Surface tension force
- Elastic compression force
- Pressure force

#### The relevant **force ratios** are:

- Froude number F = (inertial force/gravity force)<sup>1/2</sup>
- Reynolds number R = inertial force/viscous force
- Weber number W = inertial force/surface tension force
- Cauchy number C = inertial force/elastic force
- Euler number E = pressure force/inertial force

**Challenge**: Only the most relevant force ratio can be identical between a model and its prototype, if the same fluid is used, and mechanical similarity is impossible. The most relevant force ratio is selected and the remaining result in **scale effects**.

# Froude similarity $F_M = F_P$ (basic assumption)

For phenomena where gravity and inertial forces are dominant and effect of remaining forces such as kinematic viscosity are small.

Most water phenomena are modeled after Froude, particularly free surface flows (hydraulic structures, waves, wave energy converters, fluvial hydraulics including sediment transport, etc.)



Gebidem dam

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Hydraulic jump modelled after Froude

# Froude scaling laws

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| Parameter           | Dimension            | Froude          |  |  |
|---------------------|----------------------|-----------------|--|--|
|                     | Geometric similarity |                 |  |  |
| Length              | [L]                  | λ               |  |  |
| Area                | $[L^2]$              | $\lambda^2$     |  |  |
| Volume              | $[L^3]$              | $\lambda^3$     |  |  |
| Rotation            | [-]                  | 1               |  |  |
|                     | Kinematic similarity |                 |  |  |
| Time                | [T]                  | $\lambda^{1/2}$ |  |  |
| Velocity            | $[LT^{-1}]$          | $\lambda^{1/2}$ |  |  |
| Acceleration        | $[LT^{-2}]$          | 1               |  |  |
| Discharge           | $[L^{3}T^{-1}]$      | $\lambda^{5/2}$ |  |  |
|                     | Dynamic similarity   |                 |  |  |
| Mass                | [M]                  | $\lambda^3$     |  |  |
| Force               | $[MLT^{-2}]$         | $\lambda^3$     |  |  |
| Pressure and stress | $[ML^{-1}T^{-2}]$    | λ               |  |  |
| Energy and work     | $[ML^2L^{-2}]$       | $\lambda^4$     |  |  |
| Power               | $[ML^{2}T^{-3}]$     | $\lambda^{7/2}$ |  |  |

**Example** with discharge rate Q:

Given:

 $Q_M = 0.01 \text{ m}^3/\text{s}$  in a 1: $\lambda = 1:20 \text{ scale model}$ 

Unknown:

discharge in the prototype  $Q_P$ 

#### Solution:

$$Q_P = \lambda^{5/2} Q_M = 20^{5/2} 0.01 = 17.89 \text{ m}^3/\text{s}$$

# Why do scale effects matter?

Failure of Sines breakwater in 1978/9

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Strong enough in the scale model investigation (one reason for the failure were scale effects due to the incorrect scaling of the structural properties)



Images of destroyed Sines breakwater, 1978/9

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# **Illustration scale effects: Landslide-tsunamis**

Subaerial landslide-tsunami generation at 1:2 and 1:4 (Heller et al., 2008)



F = Froude number; W = Weber number; R = Reynolds number

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# **Illustration scale effects: Landslide-tsunamis**

Wave profiles of subaerial landslide-tsunamis at 1:1, 1:2 and 1:4



Scale effects relative to the maximum wave amplitude are negligible (<2%) if:</th>Reynolds number: $R \ge 300000$ Weber number: $W \ge 5000$ 

# **Avoidance: With rules of thumb**

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Satisfy limiting criteria (limiting size of the experiment) In Froude models:  $R = LV/v > R_{limit}$ ,  $W = \rho V^2 L/\sigma > W_{limit}$  etc.

In practice rules of thumb are often applied. Some examples:

- Linear wave propagation is affected less than 1% by surface tension if the wave period > 0.35 s (corresponding to a wave lenght > 0.17 m, Hughes 1993)
- Free surface water flows should be > 5 cm to avoid significant surface tension scale effects (e.g. Heller 2011)
- Wave height to measure wave force on slope during wave breaking should be larger than 0.50 m (Skladnev and Popov 1969)
- Free surface air-water flows should be conducted at W<sup>0.5</sup> > 140 (Pfister and Chanson 2012)

# **Avoidance: R invariance**

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For Froude models when **both** F and R are a priori **relevant** Horizontal axis tidal turbine



Reynolds number R = LV/v L = Length V = Velocity v = kinematic fluid viscosityAsymptotically approached power coefficient  $c_P$  level with R for a tidal energy converter suggesting limiting R = 800,000 (Bachant and Wosnik 2016)

#### Solute transport in a chlorine contact tank





 $\times 10^{6}$ 

Variation of curve area discrepancy index with scale and discharge requiring a certain model size  $\lambda \le 24$ corresponding to a limiting R (Teixeira and Rauen 2014)

# **Avoidance: Replacement of fluid**

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Water replaced by water-isopropyl alcohol (reduced surface tension, increased W)



Stagonas et al. (2011)

Water replaced by air (reduced kinematic viscosity, increased R)





### Compensation

Compensation is achieved by distorting a model geometry by giving up exact geometric similarity of some parameters in favour of an improved model-prototype similarity.

#### **Examples:**

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- Distorted models: the length λ<sub>L</sub> scale factor of a model (typically of a river) is smaller than the height and width scale factor λ to compensate increased viscose and surface tension effects with a larger flow velocity
- The grain diameter  $d_g$  in sediment transport is often not scaled with the same scale factor  $\lambda$  as the model main dimensions to avoid  $d_g < 0.5$  mm for which cohesion becomes relevant. A larger grain diameter is selected and the grain density is reduced to compensate for this



Distorted river model



Sand hold together by cohesion 15





# Work of PhD student Daniele Catucci



- He derived novel scaling laws (NSL) analytically from the governing equations (including surface tension) with the **Lie group transformation**
- He proves then numerically with OpenFOAM that these ideal scaling laws involve **no scale effects**

**Continuity Equation** 

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$$\frac{\partial U_j}{\partial x_j} = 0$$

One parameter Lie-group transformation

$$\phi = \beta^{\alpha_{\phi}} \bar{\phi}$$

**Reynolds-averaged Navier-Stokes Equation** 

$$\frac{\partial U_{i}}{\partial t} + U_{j}\frac{\partial U_{i}}{\partial x_{j}} = \frac{\partial}{\partial x_{j}}\left(\nu\frac{\partial U_{i}}{\partial x_{j}} - \overline{u_{i}u_{j}}\right) - \frac{1}{\rho}\frac{\partial p}{\partial x_{i}} + g_{i} + \frac{1}{\rho}f_{\sigma}$$

 $\alpha_{\phi}$  = Scaling exponent

- $\beta =$  Scaling parameter
- $\phi$  = Variable in the original space
- $\overline{\phi}$  = Variable in the scaled space

Surface tension

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### **Work of PhD student Daniele Catucci**

|  | scaling conditions in terms of $\alpha_x$ , $\alpha_t$ and $\alpha_ ho$<br>(novel scaling laws) |   | scaling conditions in terms of $lpha_x$ ,<br>$lpha_ ho=lpha_g=$ 0 (precise Froude scaling laws) |                                       |
|--|---|---|---|---------------------------------------|
| variables  | exponents   | scaling ratios                            | exponents   | scaling ratios                        |
| length (m)   | $\alpha_{x}$  | $\beta^{lpha_{\chi}}$                     | $\alpha_{x}$  | $\beta^{lpha_{\chi}} = \lambda$       |
| time (s)   | $\alpha_t$  | $\beta^{lpha_t}$                          | $\alpha_t = 0.5 \alpha_x$   | $\beta^{0.5\alpha_x} = \lambda^{0.5}$ |
| density (kg m $^{-3}$ )  | $\alpha_{ ho}$  | $eta^{lpha_ ho}$                          | $\alpha_{ ho}=0$  | $\beta^0 = 1$                         |
| velocity (m s $^{-1}$ )  | $\alpha_U = \alpha_x - \alpha_t$  | $\beta^{\alpha_x-\alpha_t}$               | $\alpha_U = 0.5 \alpha_x$   | $\beta^{0.5\alpha_x} = \lambda^{0.5}$ |
| pressure (Pa)  | $\alpha_p = 2\alpha_x - 2\alpha_t + \alpha_\rho$  | $\beta^{2\alpha_x-2\alpha_t+\alpha_ ho}$  | $\alpha_p = \alpha_x$   | $\beta^{lpha_{x}} = \lambda$          |
| gravitational acceleration (m s <sup>-2</sup> )  | $\alpha_g = \alpha_x - 2\alpha_t$   | $\beta^{\alpha_x-2\alpha_t}$              | $\alpha_g = 0$  | $\beta^0 = 1$                         |
| viscosity (m <sup>2</sup> s <sup>-1</sup> )  | $\alpha_{\nu}=2\alpha_{x}-\alpha_{t}$   | $\beta^{2\alpha_x-\alpha_t}$              | $\alpha_{\nu} = 1.5 \alpha_{\chi}$  | $\beta^{1.5\alpha_x} = \lambda^{1.5}$ |
| surface tension (N $\mathrm{m}^{-1}$ )   | $\alpha_{\sigma} = 3\alpha_{\chi} - 2\alpha_t + \alpha_{\rho}$                                  | $\beta^{3\alpha_x-2\alpha_t+\alpha_\rho}$ | $\alpha_{\sigma}=2\alpha_{x}$   | $\beta^{2\alpha_x} = \lambda^2$       |
| curvature of the free surface (1 $m^{-1}$ )  | $\alpha_{\kappa} = \alpha_{\chi}^{-1}$  | $\beta^{\alpha_x^{-1}}$                   | $\alpha_{\kappa} = \alpha_{\chi}^{-1}$  | $\beta^{\alpha^{-1}} = \lambda^{-1}$  |
| eddy viscosity (m <sup>2</sup> s <sup>-1</sup> )   | $\alpha_{\nu_t} = 2\alpha_x - \alpha_t$   | $\beta^{2\alpha_x-\alpha_t}$              | $\alpha_{\nu_t} = 1.5 \alpha_x$   | $\beta^{1.5\alpha_x} = \lambda^{1.5}$ |
| Reynolds stresses (m² s <sup>-2</sup> )  | $\alpha_{\langle u_i,u_j\rangle}=2\alpha_x-2\alpha_t$   | $\beta^{2\alpha_x-2\alpha_t}$             | $\alpha_{\langle u_i,u_j\rangle} = \alpha_x$  | $\beta^{lpha_{x}}=\lambda$            |
| turbulent kinetic energy (m² s <sup>-2</sup> )   | $\alpha_k = 2\alpha_x - 2\alpha_t$  | $\beta^{2\alpha_x-2\alpha_t}$             | $\alpha_k = \alpha_x$   | $\beta^{lpha_{\chi}} = \lambda$       |
| dissipation (m <sup>2</sup> s <sup><math>-3</math></sup> )                                     | $\alpha_{\epsilon}=2\alpha_{x}-3\alpha_{t}$   | $\beta^{2\alpha_{\chi}-3\alpha_{t}}$      | $\alpha_{\epsilon} = 0.5 \alpha_{x}$  | $\beta^{0.5lpha_x} = \lambda^{0.5}$   |
| production of turbulence due to horizontal velocity gradients (m <sup>2</sup> s <sup>3</sup> ) | $\alpha_{P_k}=2\alpha_x-3\alpha_t$  | $eta^{2lpha_{\chi}-3lpha_{t}}$            | $\alpha_{P_k}=0.5\alpha_x$  | $\beta^{0.5\alpha_x} = \lambda^{0.5}$ |

Catucci et al. (2021)

- The Froude scaling laws are a special case of the novel scaling laws (NSLs)
- The NSLs are more universal and flexible than the Froude scaling laws 17



#### Validation NSLs with dam break wave in OpenFOAM

Past case



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### Validation NSLs with dam break wave in OpenFOAM



• The novel scaling laws involve no scale effects in contrast to Froude scaling

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### Validation NSLs with dam break wave in OpenFOAM

Investigation of Weber number effects with phase fraction  $\gamma$ 



 The curves based on the novel scaling laws collapse (no scale effects); the ones based on Froude scaling result in scale effects 20 University of Nottingham

### Validation NSLs with dam break wave in OpenFOAM

Investigation of Reynolds number effects with turbulent kinetic energy



 The curves based on the novel scaling laws collapse (no scale effects); the ones based on Froude scaling result in scale effects



# Outlook

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- Ganular slides (Drs Sazeda Begam & Matthew Kesseler)
- Validate NSLs in the lab and NSLs including air compressibility





Side view of slide deposits from experiments at 1:1, 1:2 and 1:4 (Kesseler et al., 2020)



# Outlook

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• Wave impact on flexible structures (Tommaso Attili)

#### Time: 0.02 s





# Conclusions

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- The basics of **Froude similarity**, applied for nearly 100 years in hydraulic modelling, has been covered
- A major limitation of Froude scaling laws are **scale effects**, which become more dominant with decrease model size
- Strategies to deal with significant scale effects have been introduced, including **avoidance and compensation**
- Novel scaling laws, which are more flexible and universal that the Froude scaling laws, have been derived to model **air-water flows**
- **These novel scaling laws** work theoretically and were validated with numerical experiments, but they still need to be validated within laboratory experiments
- A range of **ongoing research projects** to understand and achieve similitude in hydraulic modelling have been shown



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